Neutrino Mass Matrices with Two Vanishing Elements/Cofactors

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Abstract

We study the phenomenological implications of the recent neutrino data for class B of two texture zeros and two vanishing cofactors for Majorana neutrinos in the flavor basis. We find that classes $B_1(B_2)$ of two texture zeros and classes $B_5(B_6)$ of two vanishing cofactors have similar predictions for neutrino oscillation parameters for the same mass hierarchy. Similar predictions for classes $B_3(B_4)$ of two texture zeros and classes $B_3(B_4)$ of two vanishing cofactors are expected. However, a preference for a shift in the quadrant of the Dirac-type CP violating phase(δ) in contrast to the earlier analysis has been predicted for a relatively large value of the reactor neutrino mixing angle(θ_{13}) for class B of two texture zeros and two vanishing cofactors for an inverted mass spectrum. No such shift in the quadrant of δ has been found for the normal mass spectrum.

1 Introduction

There has been a significant progress in precise determination of neutrino masses and mixings during the recent past. Recently, a non-zero and relatively large value of the last unknown mixing angle (θ_{13}) has been confirmed by a number of neutrino oscillation experiments [1–5] with a best fit around $9°$. Relatively large value of θ_{13} has brightened the prospects for exploring the Dirac-type CP-violating phase δ in the lepton sector, determining the octant of atmospheric mixing angle (θ_{23}) and identifying neutrino mass hierarchy. However, the currently available data on neutrino masses and mixings is not enough to determine the neutrino mass matrix(M_ν) unambiguously and additional theoretical inputs such as zero textures [6–8], vanishing minors [9–11], hybrid textures [12] and equalities between elements [13] amongst others which restrict the structure of M_ν become necessary for further progress.

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Neutrino mass matrices with two texture zeros or two vanishing cofactors are the most extensively studied texture structures. In the flavour basis, with the charged lepton mass matrix diagonal, there are 15 possible cases of two texture zeros (TZ) in the effective neutrino mass matrix M_ν of which only 7 are compatible with the current neutrino oscillation data [14]. Similarly, in case of two vanishing cofactors (TC) only 7 out of the possible 15 texture structures survive the current neutrino experimental constraints. The viable texture structures in case of TZ and TC are listed in Table 1. A_1, A_2, B_3 and B_4 texture structures of TZ in the nomenclature of Ref. [6] are equivalent to A_1, A_2, B_3 and B_4 class of TC, respectively in the nomenclature of Ref. [9]. B_1, B_2 and C of texture zeros do not appear as two zero cofactors and B_5, B_6 and D of TC provide non-trivial vanishing cofactors when confronted with the neutrino oscillation data.

Class	Texture Zeros	Class	Zero Cofactors
A_1	M_{11}, M_{12}	A_1	C_{33}, C_{32}
A_2	M_{11}, M_{13}	A_2	C_{22}, C_{32}
B_1	M_{13}, M_{22}	B_3	C_{33}, C_{13}
B_2	M_{12}, M_{33}	B_4	C_{22}, C_{21}
B_3	M_{12}, M_{22}	B_5	C_{33}, C_{12}
B_4	M_{13}, M_{33}	B_6	C_{22}, C_{13}
C	M_{22}, M_{33}	D	C_{33}, C_{22}

Table 1: Experimentally viable classes of two texture zeros [6](two vanishing cofactors [9]), here $M_{ij}(C_{ij})$ denotes the texture zero(vanishing cofactors) for the $(ij)^{th}$ element of M_ν.

Even before the confirmation of relatively large value of θ_{13}, TZ and TC have been studied comprehensively in the recent past. The signatures for distinguishing the various texture structures TZ and TC on the basis of octant of θ_{23}, quadrant of δ and mass hierarchy have been given in Ref. [7,9]. After the confirmation of relatively large value of θ_{13}, these texture structures have been further checked for their viability and predictions for the unknown neutrino parameters like δ, Majorana phases and neutrino mass scale [15]. The indistinguishability of TZ and TC has been studied in [9]. Recently, Liao et al. [16] have discussed the indistinguishability of $B_1(B_2)$ class of TZ and $B_6(B_5)$ class of TC with opposite mass hierarchy, respectively. Their analysis claims similar predictions for the oscillation parameters and Majorana phases which means that if mass hierarchy is not known texture $B_1(B_2)$ of TZ and $B_6(B_5)$ of TC are indistinguishable but the results obtained earlier in [7,9] seem to distinguish the TZ and TC on the basis of quadrant of δ. The recent global neutrino analyses favour δ around $3\pi/2$ [17,18] which motivates us to re-analyse the TZ and TC for their indistinguishability and predictions for the quadrant of δ with the recent neutrino experimental data. Texture $B_3(B_4)$ of TZ appears as $B_3(B_4)$ of TC and, therefore, shows identical predictions. Classes (B_1, B_2) of TZ and classes (B_5, B_6) of TC are P_{23}[1] symmetric texture structures, respectively. Therefore, these texture structures in TZ or in TC differ only in their predictions for the octant of θ_{23} and quadrant of δ. A close look at $B_3(B_4)$ of TZ and $B_3(B_4)$ of TC reveals that zeros and cofactors in these textures are 2-3 interchange symmetric. Similar interchange symmetry is present in $B_1(B_2)$ of TZ and $B_5(B_6)$ of TC structures. Unlike $B_3(B_4)$ of TZ and TC having trivial texture structure, it is interesting to see how the $B_1(B_2)$ of TZ is related to $B_5(B_6)$ of TC for similar mass spectrum.

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[1]For P_{23} symmetry see equations[18-20]

In this work, we present a phenomenological analysis for class B of TZ and TC in the light of the recent experimental data. The identical results for $B_3(B_4)$ of TZ and $B_3(B_4)$ of TC are trivial. However, we observe an interesting pattern of correlation between neutrino oscillation parameters for $B_1(B_2)$ of TZ and $B_5(B_6)$ of TC. Also, a relatively large value of θ_{13} has interesting implications for the quadrant of Dirac-type CP-violating phase δ. A particular quadrant for δ has been predicted for normal mass spectrum as predicted earlier in Refs. [7,9]. However, a curious preference for a shift in the quadrant of δ in contrast to the results reported in Refs. [7,9] has been predicted for an inverted mass spectrum for texture structures in class B of TZ and TC.

2 Formalism

In the flavor basis, where the charged lepton mass matrix M_l is diagonal, the complex symmetric Majorana neutrino mass matrix can be diagonalized by a unitary matrix V' as

$$M_\nu = V' M_\nu^{diag} V'^T \tag{1}$$

where $M_\nu^{diag} = \text{diag}(m_1, m_2, m_3)$. The unitary matrix V' can be parametrized as

$$V' = P_l V \quad \text{with} \quad V = U P_\nu \tag{2}$$

where [20]

$$U = \begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta} & s_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta} & c_{23}c_{13} \end{pmatrix} \tag{3}$$

with $s_{ij} = \sin\theta_{ij}$ and $c_{ij} = \cos\theta_{ij}$ and

$$P_\nu = \begin{pmatrix} 1 & 0 & 0 \\ 0 & e^{i\alpha} & 0 \\ 0 & 0 & e^{i(\beta+\delta)} \end{pmatrix}, \quad P_l = \begin{pmatrix} e^{i\varphi_e} & 0 & 0 \\ 0 & e^{i\varphi_\mu} & 0 \\ 0 & 0 & e^{i\varphi_\tau} \end{pmatrix}.$$

P_ν is the diagonal phase matrix containing two Majorana-type CP- violating phases α, β and one Dirac-type CP-violating phase δ. The phase matrix P_l is physically unobservable and depends on the phase convention. The matrix V is called the neutrino mixing matrix or the Pontecorvo-Maki-Nakagawa-Sakata (PMNS) matrix [21]. Using Eqs.(1) and (2), the effective Majorana neutrino mass matrix can be written as

$$M_\nu = P_l U P_\nu M_\nu^{diag} P_\nu^T U^T P_l^T. \tag{4}$$

The Dirac-type CP violation in neutrino oscillation experiments can be described through a rephasing invariant quantity, J_{CP} [22] with $J_{CP} = Im(U_{e1}U_{\mu2}U_{e2}^*U_{\mu1}^*)$. In the above parametrization, J_{CP} is given by

$$J_{CP} = s_{12}s_{23}s_{13}c_{12}c_{23}c_{13}^2 \sin\delta . \tag{5}$$

2.1 Two Texture Zeros in M_ν

The simultaneous existence of two zeros in the neutrino mass matrix at (a,b) and (c,d) positions implies

$$M_{\nu(ab)} = 0 \text{ and } M_{\nu(cd)} = 0 . \tag{6}$$

The above complex equations can be written as

$$m_1 A_1 + m_2 A_2 e^{2i\alpha} + m_3 A_3 e^{2i(\beta+\delta)} = 0 , \tag{7}$$
$$m_1 B_1 + m_2 B_2 e^{2i\alpha} + m_3 B_3 e^{2i(\beta+\delta)} = 0 , \tag{8}$$

where

$$A_i = U_{ai} U_{bi} \text{ and } B_i = U_{ci} U_{di} , \tag{9}$$

with $i = 1, 2, 3$ and a,b,c,d can take values e, μ, τ. These two complex Eqs. (7) and (8) involve nine physical parameters viz. the three neutrino masses (m_1, m_2, m_3), three mixing angles (θ_{12}, θ_{23}, θ_{13}), two Majorana-type CP-violating phases (α, β) and one Dirac-type CP-violating phase (δ). The masses m_2 and m_3 can be calculated from the mass-squared differences Δm_{21}^2 and $|\Delta m_{32}^2|$ using the relations

$$m_2 = \sqrt{m_1^2 + \Delta m_{21}^2} , \quad m_3 = \sqrt{m_2^2 + |\Delta m_{32}^2|} \tag{10}$$

where $\Delta m_{ij}^2 = m_i^2 - m_j^2$, $m_2 < m_3$ for Normal Hierarchy (NH) and $m_2 > m_3$ for Inverted Hierarchy (IH). Simultaneously solving Eqs. (7) and (8) for the two mass ratios, we obtain

$$\frac{m_1}{m_2} e^{-2i\alpha} = \frac{A_2 B_3 - A_3 B_2}{A_3 B_1 - A_1 B_3} \tag{11}$$

and

$$\frac{m_1}{m_3} e^{-2i\beta} = \frac{A_3 B_2 - A_2 B_3}{A_2 B_1 - A_1 B_2} e^{2i\delta} . \tag{12}$$

The magnitudes of the two mass ratios in Eqs.(11) and (12) are given by

$$\rho = \left| \frac{m_1}{m_3} e^{-2i\beta} \right| , \tag{13}$$

$$\sigma = \left| \frac{m_1}{m_2} e^{-2i\alpha} \right| . \tag{14}$$

while the CP- violating Majorana phases α and β are given by

$$\alpha = -\frac{1}{2} arg \left(\frac{A_2 B_3 - A_3 B_2}{A_3 B_1 - A_1 B_3} \right) , \tag{15}$$

$$\beta = -\frac{1}{2} arg \left(\frac{A_3 B_2 - A_2 B_3}{A_2 B_1 - A_1 B_2} e^{2i\delta} \right) . \tag{16}$$

Using the experimental inputs of the two mass-squared differences and the three mixing angles we can constrain the other parameters. Since Δm_{21}^2 and $|\Delta m_{23}^2|$ are known experimentally, the

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values of mass ratios (ρ, σ) from Eqs. (13) and (14) can be used to calculate m_1. This can be done by using Eq. (10) to obtain the two values of m_1 viz.,

$$m_1 = \sigma \sqrt{\frac{\Delta m_{21}^2}{1 - \sigma^2}} \ , \quad m_1 = \rho \sqrt{\frac{\Delta m_{21}^2 + |\Delta m_{32}^2|}{1 - \rho^2}} \ . \tag{17}$$

These two values of m_1 contain the constraints of two texture zeros in terms of ρ and σ and must be equal to within the current experimental precision. There exists a permutation symmetry between different patterns [23] of two texture zeros. The corresponding permutation matrix has the following form:

$$P_{23} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix}. \tag{18}$$

The two texture zero neutrino mass matrices, therefore, are related to each other as

$$M_\nu^X = P_{23} M_\nu^Y P_{23}^T, \tag{19}$$

leading to the following relations between the neutrino oscillation parameters:

$$\theta_{12}^X = \theta_{12}^Y, \ \theta_{13}^X = \theta_{13}^Y, \ \theta_{23}^X = \frac{\pi}{2} - \theta_{23}^Y, \ \delta^X = \delta^Y - \pi \tag{20}$$

where X and Y denote the neutrino mass matrices related by the permutation symmetry P_{23}. The texture structures related by the 2-3 permutation symmetry in class B are

$$B_1 \leftrightarrow B_2, \ B_3 \leftrightarrow B_4$$

2.2 Two Vanishing Cofactors in M_ν

The simultaneous existence of two vanishing cofactors in the neutrino mass matrix implies

$$M_{\nu(pq)} M_{\nu(rs)} - M_{\nu(tu)} M_{\nu(vw)} = 0 \ , \tag{21}$$

$$M_{\nu(p'q')} M_{\nu(r's')} - M_{\nu(t'u')} M_{\nu(v'w')} = 0 \ . \tag{22}$$

These two conditions yield two complex equations viz.

$$\sum_{k,l=1}^{3} (V_{pk} V_{qk} V_{rl} V_{sl} - V_{tk} V_{uk} V_{vl} V_{wl}) m_k m_l = 0 \ , \tag{23}$$

$$\sum_{k,l=1}^{3} (V_{p'k} V_{q'k} V_{r'l} V_{s'l} - V_{t'k} V_{u'k} V_{v'l} V_{w'l}) m_k m_l = 0 \tag{24}$$

which can be rewritten as

$$m_1 m_2 A_3 e^{2i\alpha} + m_2 m_3 A_1 e^{2i(\alpha+\beta+\delta)} + m_3 m_1 A_2 e^{2i(\beta+\delta)} = 0 \ , \tag{25}$$

$$m_1 m_2 B_3 e^{2i\alpha} + m_2 m_3 B_1 e^{2i(\alpha+\beta+\delta)} + m_3 m_1 B_2 e^{2i(\beta+\delta)} = 0 \ , \tag{26}$$

where

$$\begin{aligned} A_h &= (U_{pk} U_{qk} U_{rl} U_{sl} - U_{tk} U_{uk} U_{vl} U_{wl}) + (k \leftrightarrow l) \ , \\ B_h &= (U_{p'k} U_{q'k} U_{r'l} U_{s'l} - U_{t'k} U_{u'k} U_{v'l} U_{w'l}) + (k \leftrightarrow l) \ , \end{aligned} \tag{27}$$

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with (h, k, l) as the cyclic permutation of (1,2,3). Simultaneously solving Eqs. (25) and (26) for the two mass ratios, we obtain

$$\frac{m_1}{m_2}e^{-2i\alpha} = \frac{A_3B_1 - A_1B_3}{A_2B_3 - A_3B_2} \tag{28}$$

and

$$\frac{m_1}{m_3}e^{-2i\beta} = \frac{A_2B_1 - A_1B_2}{A_3B_2 - A_2B_3}e^{2i\delta} . \tag{29}$$

The magnitudes of the two mass ratios in Eqs. (28) and (29) are given by

$$\rho = \left| \frac{m_1}{m_3}e^{-2i\beta} \right|, \tag{30}$$

$$\sigma = \left| \frac{m_1}{m_2}e^{-2i\alpha} \right| \tag{31}$$

while the CP- violating Majorana phases α and β are given by

$$\alpha = -\frac{1}{2}arg\left(\frac{A_3B_1 - A_1B_3}{A_2B_3 - A_3B_2} \right), \tag{32}$$

$$\beta = -\frac{1}{2}arg\left(\frac{A_2B_1 - A_1B_2}{A_3B_2 - A_2B_3}e^{2i\delta} \right). \tag{33}$$

Different classes of neutrino mass matrices with two vanishing cofactors are related by the permutation symmetry P_{23} just as in case of neutrino mass matrices with two texture zeros. The texture structures related by 2-3 permutation symmetry in class B are

$$B_3 \leftrightarrow B_4, B_5 \leftrightarrow B_6$$

3   Numerical Analysis

The present experimental constraints on neutrino parameters at 1, 2 and 3σ [14] are given in Table 2. For the numerical analysis, we generate 10^7 random points (10^8 when the number of allowed points is small). The effective Majorana mass of the electron neutrino $|M_{ee}|$ which determines the rate of neutrinoless double beta decay (NLDB) is given by

$$|M_{ee}| = |m_1c_{12}^2c_{13}^2 + m_2s_{12}^2c_{13}^2e^{2i\alpha} + m_3s_{13}^2e^{2i\beta}|. \tag{34}$$

The observation of NLDB decay would signal lepton number violation indicating physics beyond the standard model. Its observation will imply the Majorana nature of neutrinos and, in addition, provide a way to probe the neutrino mass scale. There are a large number of projects such as CUORICINO [24], CUORE [25], MAJORANA [26], SuperNEMO [27], EXO [28], GENIUS [29] which aim to achieve a sensitivity upto 0.01 eV for $|M_{ee}|$. In our numerical analysis, we take the upper limit of $|M_{ee}|$ to be 0.5 eV. In addition, the cosmological data put an upper limit on the sum of the active neutrino masses:

$$\Sigma = \sum_{i=1}^{3} m_i. \tag{35}$$

Parameter	Mean$^{(+1\sigma,+2\sigma,+3\sigma)}_{(-1\sigma,-2\sigma,-3\sigma)}$
$\Delta m^2_{21}[10^{-5}eV^2]$	$7.62^{(+0.19,+0.39,+0.58)}_{(-0.19,-0.35,-0.5)}$
$\Delta m^2_{31}[10^{-3}eV^2]$	$2.55^{(+0.06,+0.13,+0.19)}_{(-0.09,-0.19,-0.24)}$,
	$(-2.43^{(+0.09,+0.19,+0.24)}_{(-0.07,-0.15,-0.21)})$
$\sin^2\theta_{12}$	$0.32^{(+0.016,+0.03,+0.05)}_{(-0.017,-0.03,-0.05)}$
$\sin^2\theta_{23}$	$0.613^{(+0.022,+0.047,+0.067)}_{(-0.04,-0.233,-0.25)}$,
	$(0.60^{(+0.026,+0.05,+0.07)}_{(-0.031,-0.210,-0.230)})$
$\sin^2\theta_{13}$	$0.0246^{(+0.0028,+0.0056,+0.0076)}_{(-0.0029,-0.0054,-0.0084)}$,
	$(0.0250^{(+0.0026,+0.005,+0.008)}_{(-0.0027,-0.005,-0.008)})$

Table 2: Current Neutrino oscillation parameters from global fits [14]. The upper (lower) row corresponds to Normal (Inverted) Hierarchy, with $\Delta m^2_{31} > 0$ ($\Delta m^2_{31} < 0$).
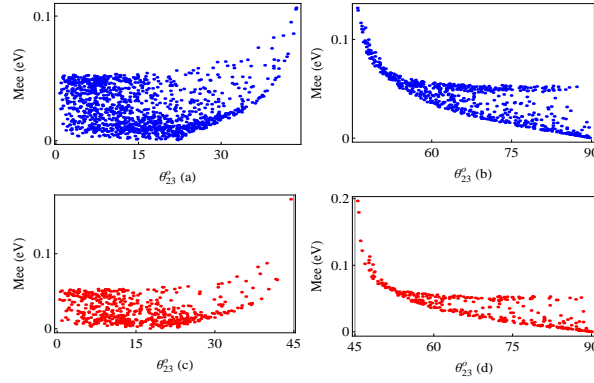

Figure 1: Correlation plots for B_3 of TZ(upper row) and TC(lower row) for NH(left panel) and IH(right panel) specturm.
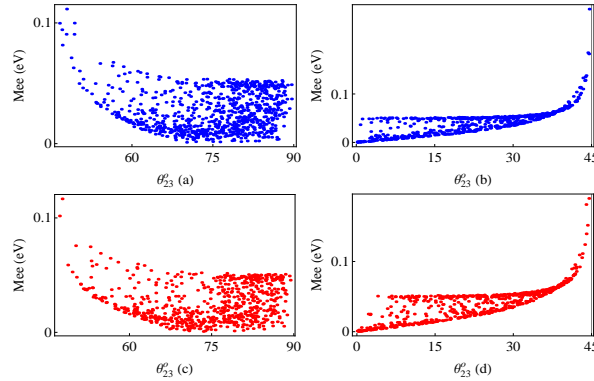

Figure 2: Correlation plots for B_4 of TZ(upper row) and TC(lower row) for NH(left panel) and IH(right panel) specturm.

The experimental results from Planck [30] limit $\sum < 0.23$ eV at 95% confidence level (CL). However, these bounds strongly depend on the model details as well as on the data set used.

In the present work, we take a conservative upper limit on $\sum$ to be $1\ eV$.

In our numerical analysis, we vary the two mass-squared differences $(\Delta m_{21}^2, \Delta m_{31}^2)$ randomly within their 3σ allowed ranges but keep the other oscillation parameters free. The three neutrino mixing angles and Dirac-type CP-violating phase δ are varied between $0°\text{-}90°$ and $0°\text{-}360°$, respectively. We perform the numerical analysis for class B of two texture zeros (TZ) and two vanishing cofactors(TC). It is observed that the octant of θ_{23} is well restricted for B_3, B_4 of TZ and TC (see Figs.[1, 2]). This result is independent of the values of solar and reactor mixing angles and no constraint on $|M_{ee}|$ is required. However, for other classes of TZ and TC the value of θ_{23} is unconstrained like other oscillation parameters.

In the second step, we use the constraint of large $|M_{ee}|$ along with the experimental range of Δm_{21}^2, Δm_{31}^2 while keeping the mixing angles and the Dirac-type CP violating phase δ free. It is found that all the textures of class B of TZ as well as TC predict near maximal neutrino atmospheric mixing θ_{23}. Using the experimental constraints on neutrino mixing angles supplemented with a large $|M_{ee}|$, class B of TZ as well as TC predict near maximal atmospheric mixing angle θ_{23} and Dirac-type CP-violating phase δ fixed near $\pi/2$ or $3\pi/2$. These results
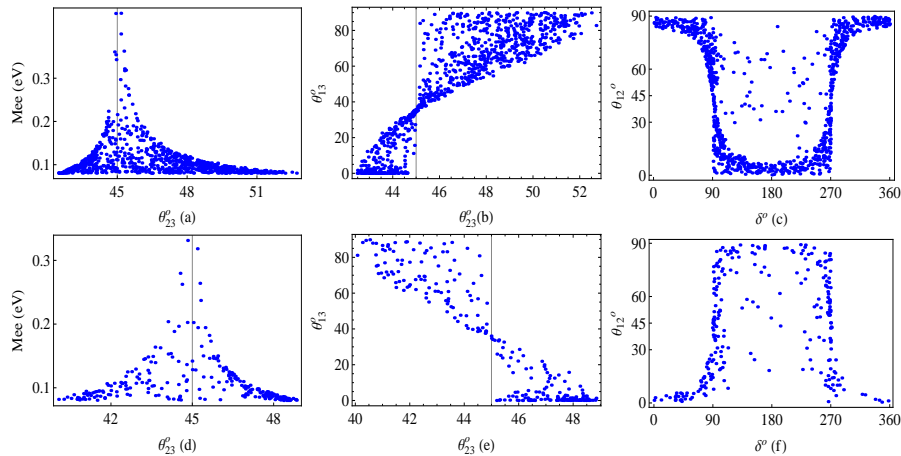

Figure 3: Correlation plots for B_1 of TZ for NH(upper row) and IH(lower row) spectrum.

are independent of the values of solar(θ_{12}) and reactor (θ_{13}) neutrino mixing angles as noted by Dev et. al. in Ref. [11] for textures B_5 and B_6 of TC and in Ref. [31] for class B of TZ. The correlation plots for the oscillation parameters are shown in Fig.[3] and Fig.[4].

It is evident from Fig.[3] and Fig.[4] that class B_1 of TZ and class B_5 of TC show similar predictions for the same mass hierarchy. It is clear from Fig.[3] and Fig.[5] that classes B_1 of TZ and B_6 of TC with opposite mass hierarchy are distinguishable on the basis of the oscillation parameter δ. This feature has not been discussed in Ref. [16] while considering the indistinguishable nature of these texture structures. This distinction, in fact, becomes more clear if texture structures under investigation are subjected to the experimental inputs of neutrino mixing angles θ_{12}, θ_{23} along with Δm_{21}^2, Δm_{31}^2 and large $|M_{ee}|$ while allowing θ_{13} to vary between $0°\text{-}12°$. All these texture structures show a marked preference for a particular quadrant of δ as given in Table 3. It is clear from the Table 3 that a distinction between $B_1(B_2)$ of TZ and $B_6(B_5)$ of TC can be made on the basis of quadrant of CP violating phase δ even if neutrino mass hierarchy is not known. It is important to keep in mind that the results obtained in Table 3 are for the values of θ_{13} between $0°$ to $12°$ as the exact value of θ_{13} was not
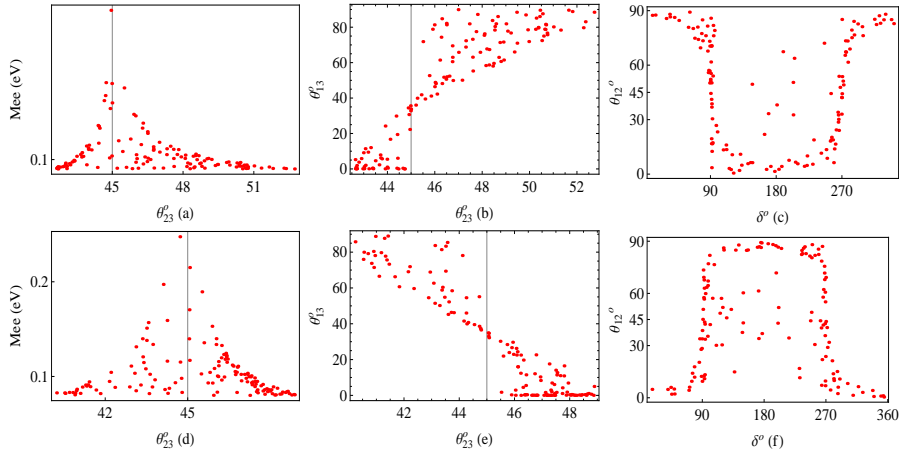
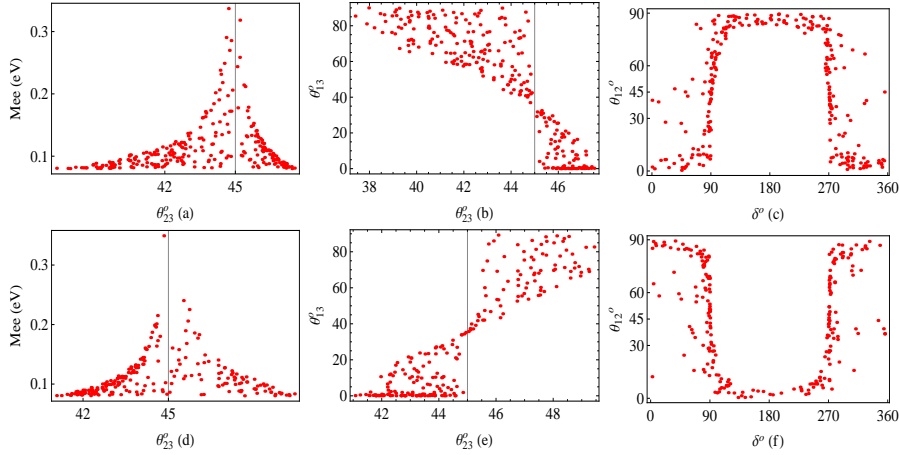Figure 4: Correlation plots for B_5 of TC for NH(upper row) and IH(lower row) spectrum.

Figure 5: Correlation plots for B_6 of TC for NH(upper row) and IH(lower row) spectrum.

known at that time. However, a change in the quadrant of δ which is in marked contrast from the results tabulated in Table 3 has been predicted for a relatively large value of θ_{13} presently known. Interestingly, this change has been predicted only for the inverted mass hierarchy while the results for the normal hierarchy spectrum are in conformity with the results from earlier analysis presented in Table 3.

In the next step, we vary δ randomly between 265°-275° and all other parameters within their 3σ CL ranges to explore the quadrant of δ more accurately in the light of the current experimental data. Fig.[6] shows correlation in B_1 of TZ and B_5 of TC for normal mass spectrum and Fig.[7] shows correlation in B_1 of TZ and B_5 of TC for an inverted mass spectrum. It is evident from the plots that classes $B_1(B_3)$ of TZ and $B_5(B_3)$ of TC show nearly similar predictions for oscillation parameters. However, Fig.[6] reveals that there is a small difference in the Majorana phase β for B_1 of TZ and B_5 of TC. Such differences can be there for $B_1(B_2)$ of TZ and $B_5(B_6)$ of TC, respectively as these textures are not related trivially as is the case for $B_3(B_4)$ of TZ and $B_3(B_4)$ of TC respectively. Distinguishing these texture structures depends on how precisely the Majorana phases are measured in the near future [32].

Class	δ quadrant	Normal Hierarchy	Inverted Hierarchy
		Two Texture Zeros	
B_1	(2,3)	$\theta_{23} < \pi/4$	$\theta_{23} > \pi/4$
B_2	(1,4)	$\theta_{23} > \pi/4$	$\theta_{23} < \pi/4$
B_3	(1,4)	$\theta_{23} > \pi/4$	$\theta_{23} < \pi/4$
B_4	(2,3)	$\theta_{23} < \pi/4$	$\theta_{23} > \pi/4$
		Two Vanishing Cofactors	
B_3	(1,4)	$\theta_{23} < \pi/4$	$\theta_{23} > \pi/4$
B_4	(2,3)	$\theta_{23} > \pi/4$	$\theta_{23} < \pi/4$
B_5	(2,3)	$\theta_{23} < \pi/4$	$\theta_{23} > \pi/4$
B_6	(1,4)	$\theta_{23} > \pi/4$	$\theta_{23} < \pi/4$

Table 3: Experimental signatures distinguishing class B of two texture zeros [7](two vanishing cofactors [9]).
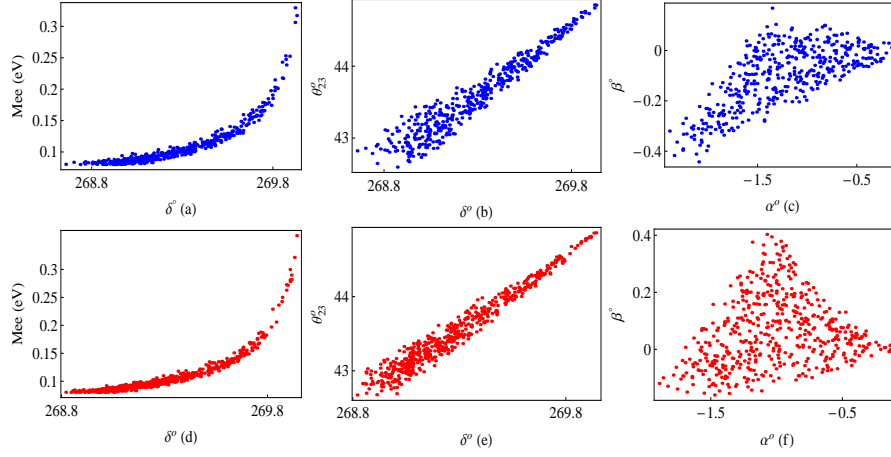

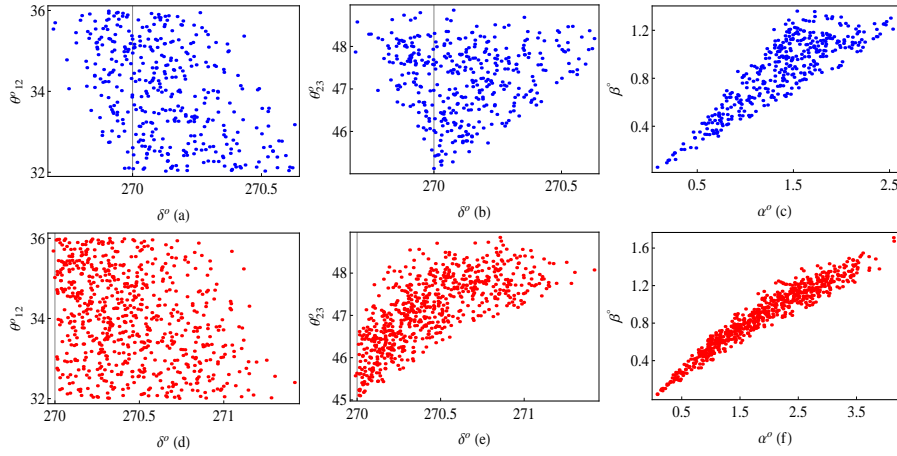Figure 6: Correlation plots for B_1 of TZ(upper row) and B_5 of TC(lower row) for NH spectrum.

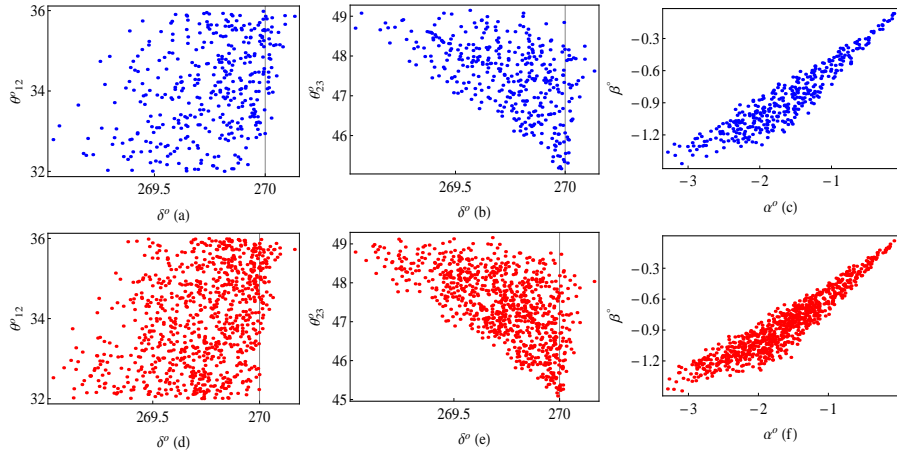Figure 7: Correlation plots for B_1 of TZ(upper row) and B_5 of TC(lower row) for IH spectrum.

Figure 8: Correlation plots for B_3 of TZ(upper row) and TC(lower row) for IH spectrum.
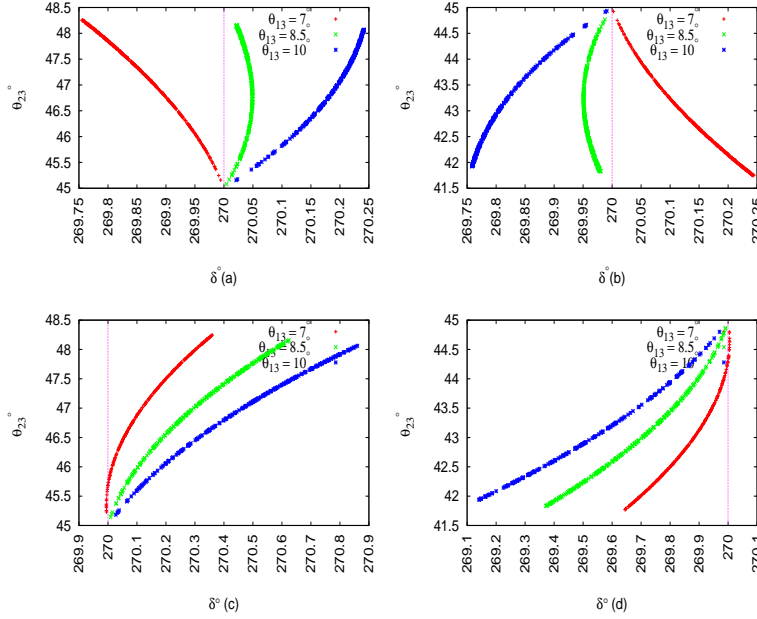

Figure 9: Correlation between θ_{23} and δ for $B_1(a)$ and $B_2(b)$ of TZ (upper row) and $B_5(c)$ and $B_6(d)$ of TC (lower row) for IH spectrum.

Fig.[8] shows correlation in textures B_3 of TZ and TC for IH spectrum. The textures B_3 of TZ and TC show similar predictions for oscillation parameters which is expected as these texture structures are trivially related to each other. The same is the case for textures B_4 of TZ and TC. Correlation plots in Fig.[1] and Fig.[2] depict the permutation symmetry between B_3 and B_4 class of TZ and TC and, also, the permutation symmetry between B_5 and B_6 texture structures of TC is evident from Fig.[4] and Fig.[5]. The results for the 2-3 interchange symmetric textures which are not discussed here can simply be obtained using Eq. (20). Fig.[9] shows the correlation plots for B_1, B_2 of TZ and B_5, B_6 of TC for IH spectrum. The plots show correlation between θ_{23} and δ for different values of θ_{13} while all other parameters are at their best fit values. It is clear from the figure that there is a preference for change in the quadrant of δ in contrast to the results tabulated in Table 3 for the presently known relatively
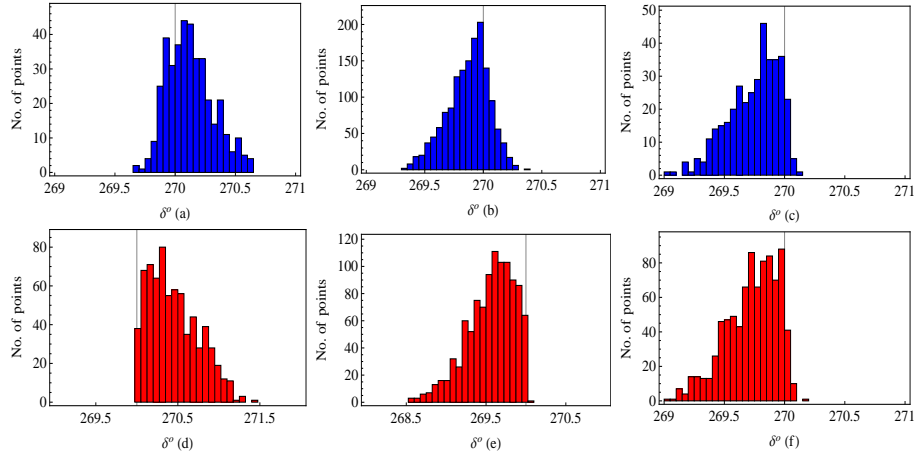
11

Figure 10: Histogram plots for B_1(a), B_2(b), B_3(c) of TZ and B_5(d), B_6(e), B_3(f) of TC for IH spectrum.

large value of reactor mixing angle θ_{13} for IH spectrum. For relatively large value of θ_{13}, the Dirac-type CP-violating phase δ shows a preference for the (1,4) quadrant in class $B_1(B_5)$ of TZ(TC) and a preference for the (2,3) quadrant in class B_3 of TZ and TC in contrast to the results summarized in Table 3. However, for normal hierarchical mass spectrum the marked preference for a particular quadrant of δ in class B of TZ and TC is still there and is the same as given in Table 3. To see the preference for shift in quadrant of δ more clearly for inverted hierarchical spectrum, histograms have been plotted for class B of TZ and TC(see Fig.[10]). The histograms clearly show the most probable values for Dirac-type CP violating phase δ varied between the range 265°-275° whereas all other neutrino oscillation parameters are varied within their experimental 3σ ranges. One can see from the histograms that all the texture structures of class B show a preference for shift in the quadrant of δ in contrast to results of the earlier analyses given in Table 3 for TZ and TC for the present neutrino oscillation data with relatively large θ_{13}.

4   Summary

In the present work, phenomenological analysis for classes B of TZ and TC has been presented in the light of the recent neutrino oscillation data. It has been noted that the quadrant of θ_{23} is independent of the values of θ_{12} and θ_{13} and does not require the constraint of large value of $|M_{ee}|$ for B_3, B_4 textures of TZ and TC. Classes $B_1(B_2)$ of TZ and $B_5(B_6)$ of TC have been found to have similar predictions for similar mass hierarchy except for a small difference in Majorana phases which can be due to the non-trivial structures of TZ and TC for these texture structures. The preference for a particular quadrant of the CP violating phase δ in contrast to the results reported in earlier analyses has been predicted for class B of TZ and TC for an inverted mass spectrum.

Acknowledgements

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